

Scattering Matrix Formulation for Substructure Characteristic Mode Analysis

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 $\begin{array}{l} \mathsf{M}. \mbox{ Gustafsson et al. "Unified theory of characteristic modes: Part I-Fundamentals". IEEE Trans. Antennas Propag. \\ 70.12 (2022), pp. 11801-11813 \\ \mathsf{M}. \mbox{ Gapek et al. "Characteristic Mode Decomposition Using the Scattering Dyadic in Arbitrary Full-Wave Solvers". \\ IEEE Trans. Antennas Propag. 71.1 (2023), pp. 830-839 \\ \end{array}$



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• Modal significance
$$|t_n| = \frac{1}{\sqrt{1+\lambda_n^2}}$$

Many equivalent representations for CM

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$$\begin{bmatrix} \mathbf{Z}_{cc} & \mathbf{Z}_{c0} \\ \mathbf{Z}_{0c} & \mathbf{Z}_{00} \end{bmatrix} \begin{bmatrix} \mathbf{I}_0 \\ \mathbf{I}_c \end{bmatrix} = \begin{bmatrix} \mathbf{0} \\ \mathbf{V}_c \end{bmatrix}$$



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 $\begin{array}{c} \Omega_{c} \\ V_{c} \\ I_{c} \\ \end{array} \begin{array}{c} \Omega_{0} \\ V_{0} = 0 \\ I_{0} \\ \end{array}$ Note: decompositions based on basis function (DoF) can require a fine mesh to model connected regions

 $\Omega_{\rm G}$

 $\Omega_{\rm A}$

Reduced MoM system matrix

$$\widetilde{\mathbf{Z}}\mathbf{I}_c = \widetilde{\mathbf{R}}\mathbf{I}_c + j\widetilde{\mathbf{X}}\mathbf{I}_c = (\mathbf{Z}_{cc} - \mathbf{Z}_{c0}\mathbf{Z}_{00}^{-1}\mathbf{Z}_{0c})\mathbf{I}_c = \mathbf{V}_c$$



J. Ethier and D. McNamara. "Sub-structure characteristic mode concept for antenna shape synthesis". Electronics letters 48.9 (2012), p. 1

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 Substructure CM from the generalized eigenvalue problem [EM12]

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What is the corresponding scattering formulation for substructure CM?



Scattering matrices



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 - ▶ S₀ for the background (uncontrollable region)



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► The scattering eigenvalues s_n are related to modal significance |t_n| and MoM substructure characteristic eigenvalues λ_n = eig(X̃, R̃) as

$$t_n = rac{s_n - 1}{2} \quad ext{and} \ \lambda_n = -\operatorname{Im}\{t_n^{-1}\} = \mathrm{j}rac{s_n + 1}{s_n - 1}$$



Scaling of the scattered field from the background





- Scaling of the scattered field from the background
- Unitary scattering matrices (lossless)

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► The scattered fields can be expressed in transition matrices (or scattering dyadics) S = 2T + 1 and S₀ = 2T₀ + 1

$$\frac{1}{2}(\mathbf{S}_{0}^{\mathrm{H}}\mathbf{S}-\mathbf{1})\mathbf{a}_{n} = (2\mathbf{T}_{0}^{\mathrm{H}}\mathbf{T}+\mathbf{T}_{0}^{\mathrm{H}}+\mathbf{T})\mathbf{a}_{n} = t_{n}\mathbf{a}_{n}$$



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Maximal scattering (difference between the scattered fields of the composite object Ta_n and the background T₀a_n)

$$|(\mathbf{T} - \mathbf{T}_0)\mathbf{a}_n|^2 = \mathbf{a}_n^{\mathrm{H}}(\mathbf{T}_0^{\mathrm{H}}\mathbf{T}_0 + \mathbf{T}^{\mathrm{H}}\mathbf{T} - 2\operatorname{Re}\{\mathbf{T}_0^{\mathrm{H}}\mathbf{T}\})\mathbf{a}_n\}$$

= - Re{ $\mathbf{a}_n^{\mathrm{H}}(2\mathbf{T}_0^{\mathrm{H}}\mathbf{T} + \mathbf{T}_0^{\mathrm{H}} + \mathbf{T})\mathbf{a}_n\}$ = - Re{ $\{t_n\}|\mathbf{a}_n|^2 = |t_n|^2|\mathbf{a}_n|^2$



• factorizing the radiation matrix $\mathbf{R} = \operatorname{Re}\{\mathbf{Z}\} = \mathbf{U}^{\mathrm{T}}\mathbf{U}$ into spherical waves [Gus+22a] with $\mathbf{U} = \begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_c \end{bmatrix}$ and setting

$$\widetilde{\mathbf{U}} = \mathbf{U}_{\mathrm{c}} - \mathbf{U}_{0}\mathbf{Z}_{00}^{-1}\mathbf{Z}_{00}$$

▶ factorizing the radiation matrix R = Re{Z} = U^TU into spherical waves [Gus+22a] with U = [U₀ U_c] and setting

 $\widetilde{\mathbf{U}} = \mathbf{U}_c - \mathbf{U}_0 \mathbf{Z}_{00}^{-1} \mathbf{Z}_{0c}$

(*)

► reformulates the CM eigenvalue problem to $\widetilde{\mathbf{Z}}\widetilde{\mathbf{I}}_n = (1 + j\lambda_n)\widetilde{\mathbf{U}}^{\mathrm{H}}\widetilde{\mathbf{U}}\widetilde{\mathbf{I}}_n \Rightarrow -\widetilde{\mathbf{U}}\widetilde{\mathbf{Z}}^{-1}\widetilde{\mathbf{U}}^{\mathrm{H}}\mathbf{f}_n = t_n\mathbf{f}_n$ where $\mathbf{f}_n = -\widetilde{\mathbf{U}}\widetilde{\mathbf{I}}_n$ and $t_n = -1/(1 + j\lambda_n)$

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$$\mathbf{T} = -\mathbf{U}\mathbf{Z}^{-1}\mathbf{U}^{\mathrm{T}} = -\begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_\mathrm{c} \end{bmatrix} \begin{bmatrix} \mathbf{Z}_{00} & \mathbf{Z}_{0\mathrm{c}} \\ \mathbf{Z}_{\mathrm{c}0} & \mathbf{Z}_{\mathrm{cc}} \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{U}_0^{\mathrm{T}} \\ \mathbf{U}_\mathrm{c}^{\mathrm{T}} \end{bmatrix} \quad \mathbf{T}_0 = -\mathbf{U}_0\mathbf{Z}_{00}^{-1}\mathbf{U}_0^{\mathrm{T}}$$

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Transition matrices of the composite object Ω and background object Ω₁ are expressed in MoM system matrices

$$\mathbf{T} = -\mathbf{U}\mathbf{Z}^{-1}\mathbf{U}^{\mathrm{T}} = -\begin{bmatrix} \mathbf{U}_0 & \mathbf{U}_c \end{bmatrix} \begin{bmatrix} \mathbf{Z}_{00} & \mathbf{Z}_{0c} \\ \mathbf{Z}_{c0} & \mathbf{Z}_{cc} \end{bmatrix}^{-1} \begin{bmatrix} \mathbf{U}_0^{\mathrm{T}} \\ \mathbf{U}_c^{\mathrm{T}} \end{bmatrix} \quad \mathbf{T}_0 = -\mathbf{U}_0\mathbf{Z}_{00}^{-1}\mathbf{U}_0^{\mathrm{T}}$$

▶ Block inversion shows that MoM matrix (*) and scattering based substructure modes $2\mathbf{T}_0^H\mathbf{T} + \mathbf{T}_0^H + \mathbf{T}$ are equivalent

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Numerical example



- (left) full structure (right) substructure
- Adapted from [EM12], $\ell = 120 \,\mathrm{mm}$, $w = 60 \,\mathrm{mm}$, $h = 15 \,\mathrm{mm}$, $d = 30 \,\mathrm{mm}$
- Negligible differences for such as a sheet resistance $0.01 \Omega/\Box$

Generalizations

Many possible generalizations:

- Embedded structures
- Cavities
- Combination with ports



Generalizations



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Scattering based formulation of substructure CM



M. Gustafsson et al. "Unified theory of characteristic modes: Part I-Fundamentals". *IEEE Trans. Antennas Propag.* 70.12 (2022), pp. 11801–11813; M. Gustafsson et al. "Unified theory of characteristic modes: Part II-Tracking, losses, and FEM evaluation". *IEEE Trans. Antennas Propag.* 70.12 (2022), pp. 11814–11824

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 Physical insight



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- Scattering based formulation of substructure CM
- Physical insight
- $\blacktriangleright\,$ MoM, FEM, or FDTD to compute matrix ${\bf S}$ and ${\bf T}$



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- \blacktriangleright MoM, FEM, or FDTD to compute matrix ${f S}$ and ${f T}$
- Favorite MoM formulation, EFIE, MFIE, CFIE, PMCHWT, Nyström



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